



Evidence for the chiral anomaly in a Dirac Semi-metal



Jun Xiong



S. Kushwaha



Tian Liang



Jason Krizan



Quinn Gibson



Bob Cava

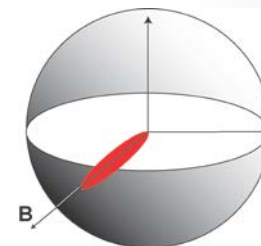


NPO

Jun Xiong, Tian Liang, Minhao Liu, Max Hirschberger, Wudi Wang, N. P. Ong
Department of Physics, Princeton Univ.

Satya Kushwaha, Quinn Gibson, Jason Krizan, Maz Ali, R. J. Cava
Department of Chemistry, Princeton Univ.

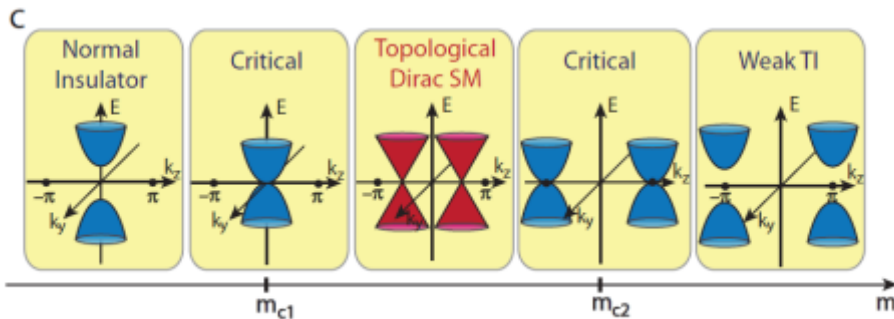
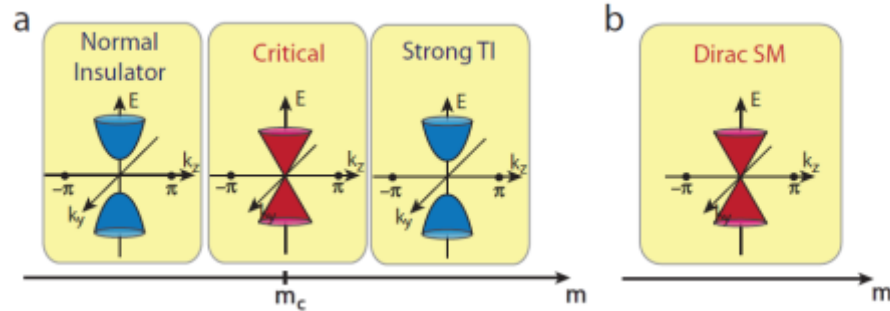
1. Introduction to Dirac semimetals
2. Ultrahigh mobilities and giant MR in Cd_3As_2
3. Intro to Weyl physics
4. Chiral anomaly in Na_3Bi



Supported by MURI, ARO, NSF (MRSEC), Moore Foundation

Dirac node protection by Point-Group Symmetry

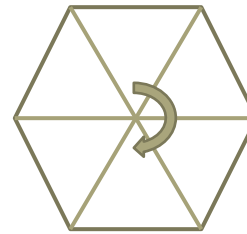
unstable



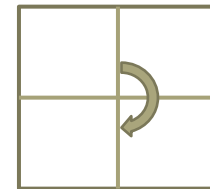
Nodes stable

Wan, Turner, Vishwanath, Savrasov, *PRB* 2011
 Young, Kane, Mele et al. *PRL* 2012
 Wang, Dai et al., *PRB* 2012
 Fang, Gilbert, Dai, Bernevig, *PRL* 2012

B.J. Yang and N. Nagaosa, arXiv
 1404.0754



C_3



C_4

Initial idea: Dirac node is protected by TRS and IS if *pinned* to zone corners.

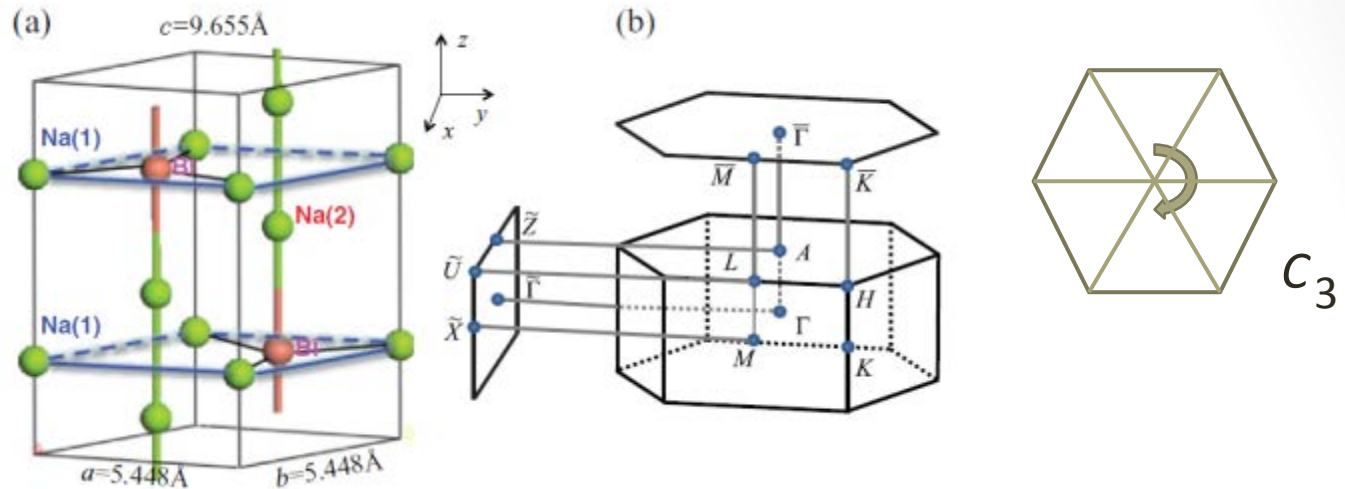
Later, inclusion of point-group symmetry C_n extends protection to *anywhere* on symmetry axis.

Protected Dirac nodes in semimetal Na_3Bi

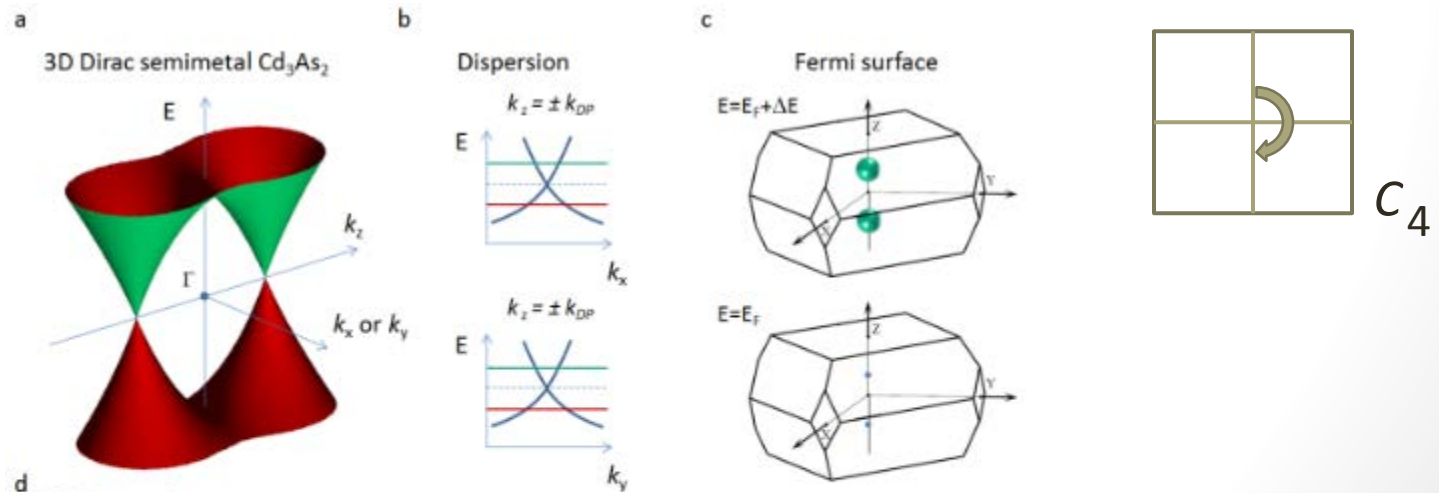
Wang, Dai, Fang et al. *PRB* 2012

Wang, Dai, Fang et al. *PRB* 2013

Na_3Bi



Cd_3As_2



Three Dimensional Dirac Semimetal and Quantum Spin Hall Effect in Cd_3As_2

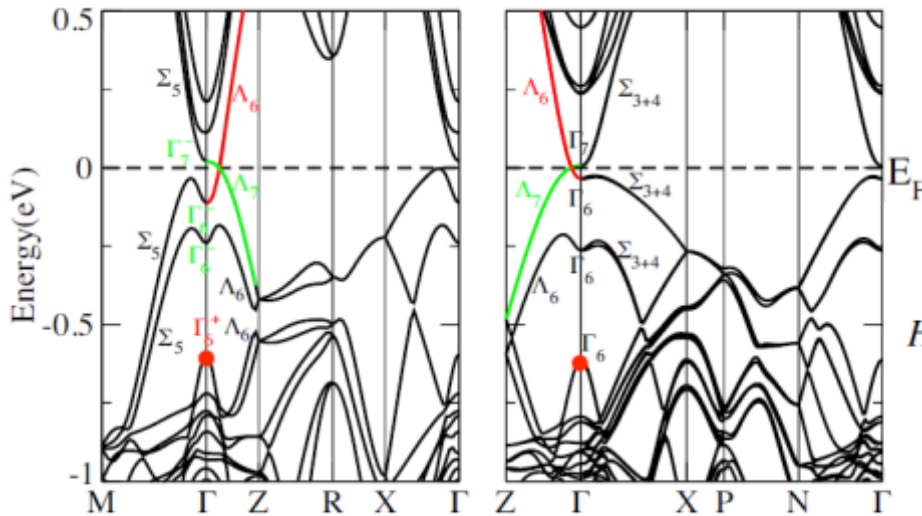
Zhijun Wang, Hongming Weng,* Quansheng Wu, Xi Dai, and Zhong Fang[†]

IOP, Beijing

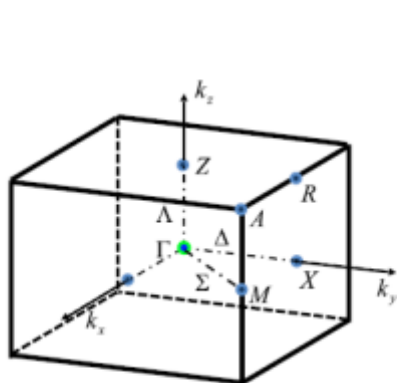
Point Group C_4

4x4 basis set

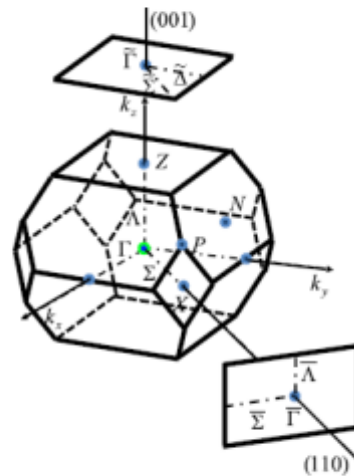
$|S, \frac{1}{2}\rangle, |P, \frac{3}{2}\rangle, |S, -\frac{1}{2}\rangle, |P, -\frac{3}{2}\rangle$



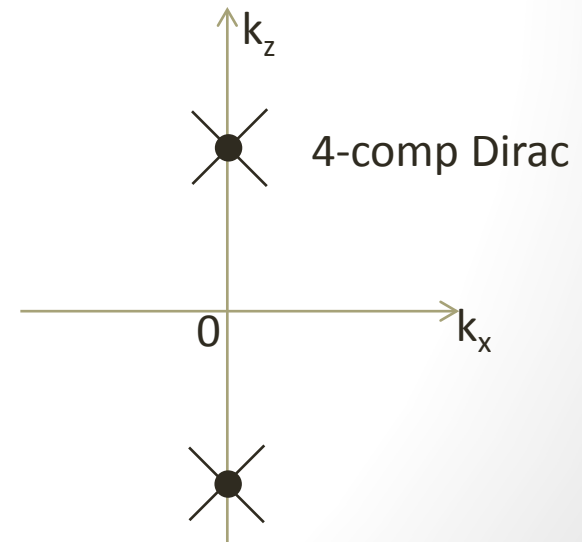
$$H_{\Gamma}(\mathbf{k}) = \epsilon_0(\mathbf{k}) + \begin{pmatrix} M(\mathbf{k}) & Ak_+ & 0 & B^*(\mathbf{k}) \\ Ak_- & -M(\mathbf{k}) & B^*(\mathbf{k}) & 0 \\ 0 & B(\mathbf{k}) & M(\mathbf{k}) & -Ak_- \\ B(\mathbf{k}) & 0 & -Ak_+ & -M(\mathbf{k}) \end{pmatrix},$$



(a) structure I

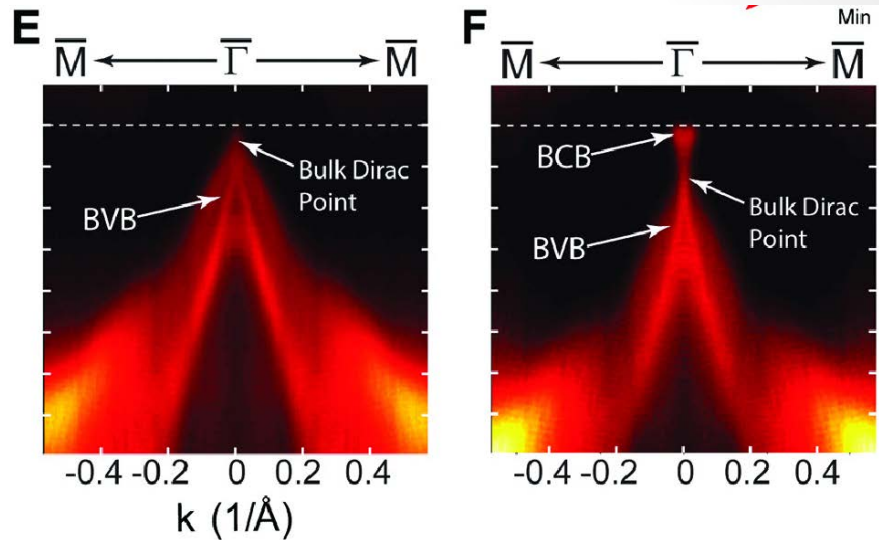
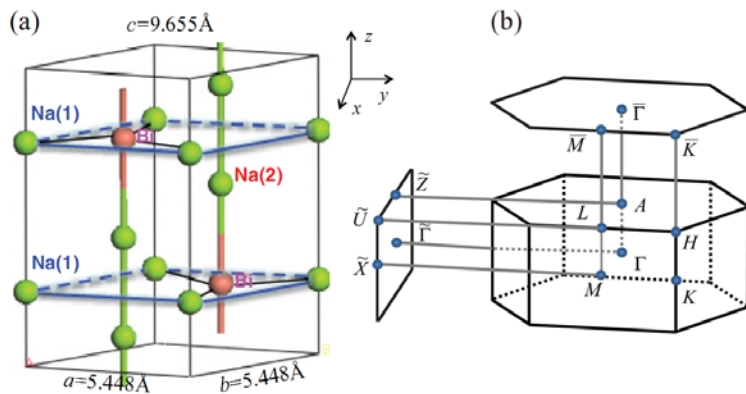


(b) structure II



Photoemission on Na_3Bi and Cd_3As_2

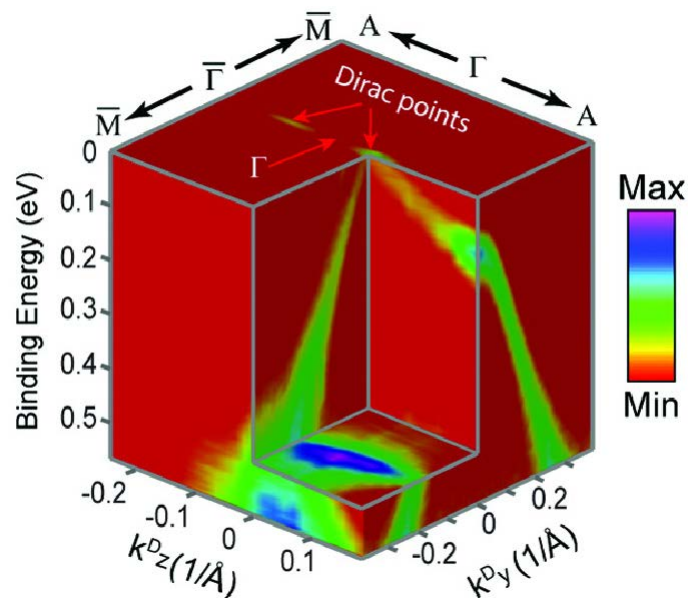
Point group C3



Liu, Chen *et al*, *Science* 2014

Neupane, Hasan *et al.*, *Nat Comm.* 2014

Borisenko, Cava *et al.*, *PRL* 2014

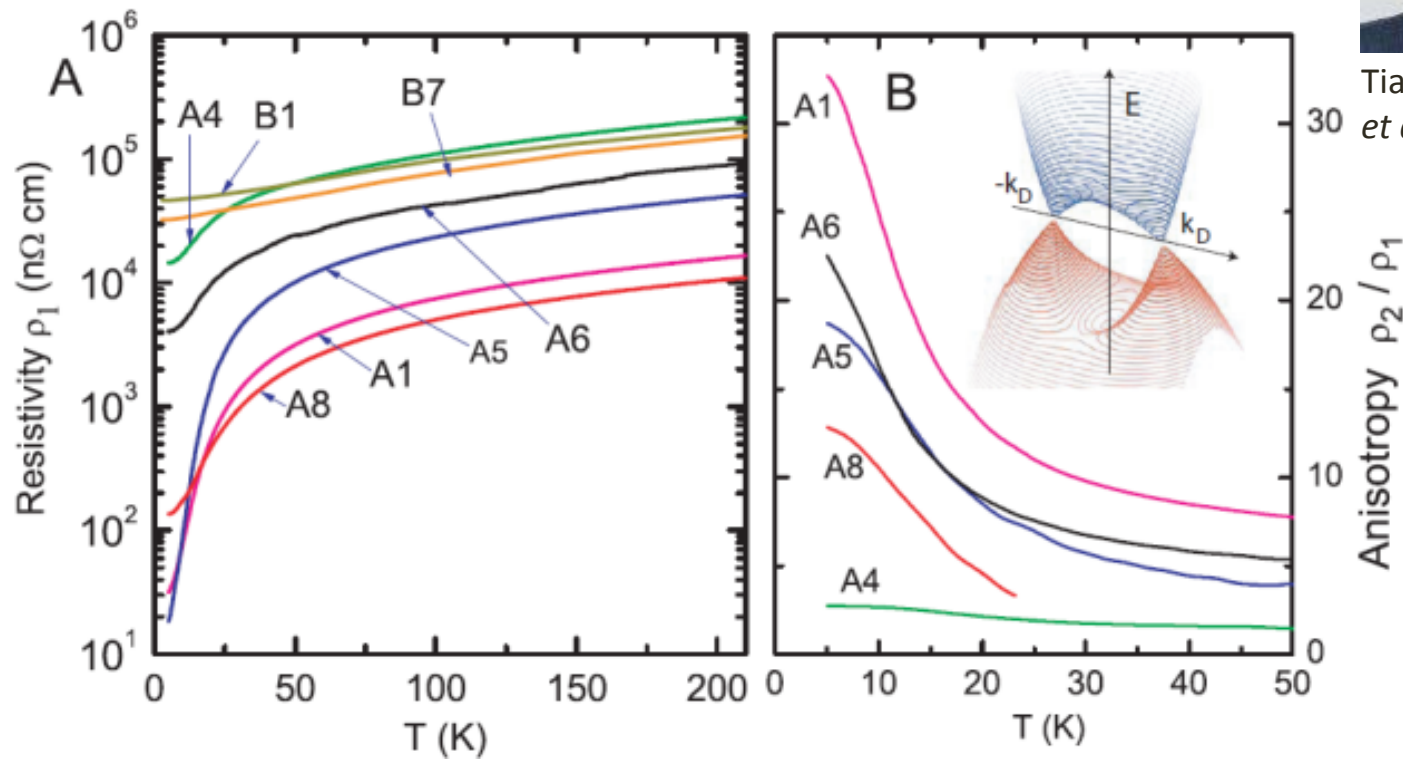


$$V_x \approx V_y = 3.74 \times 10^5 \text{ m/s}, V_z = 2.89 \times 10^4 \text{ m/s}$$

Cd₃As₂ Resistivity vs. Temperature --- RRR shows huge variation



Tian Liang, Quinn Gibson,
et al. Nat. Mater. 2015



Set A (crystals) show 1000 fold difference in RRR.

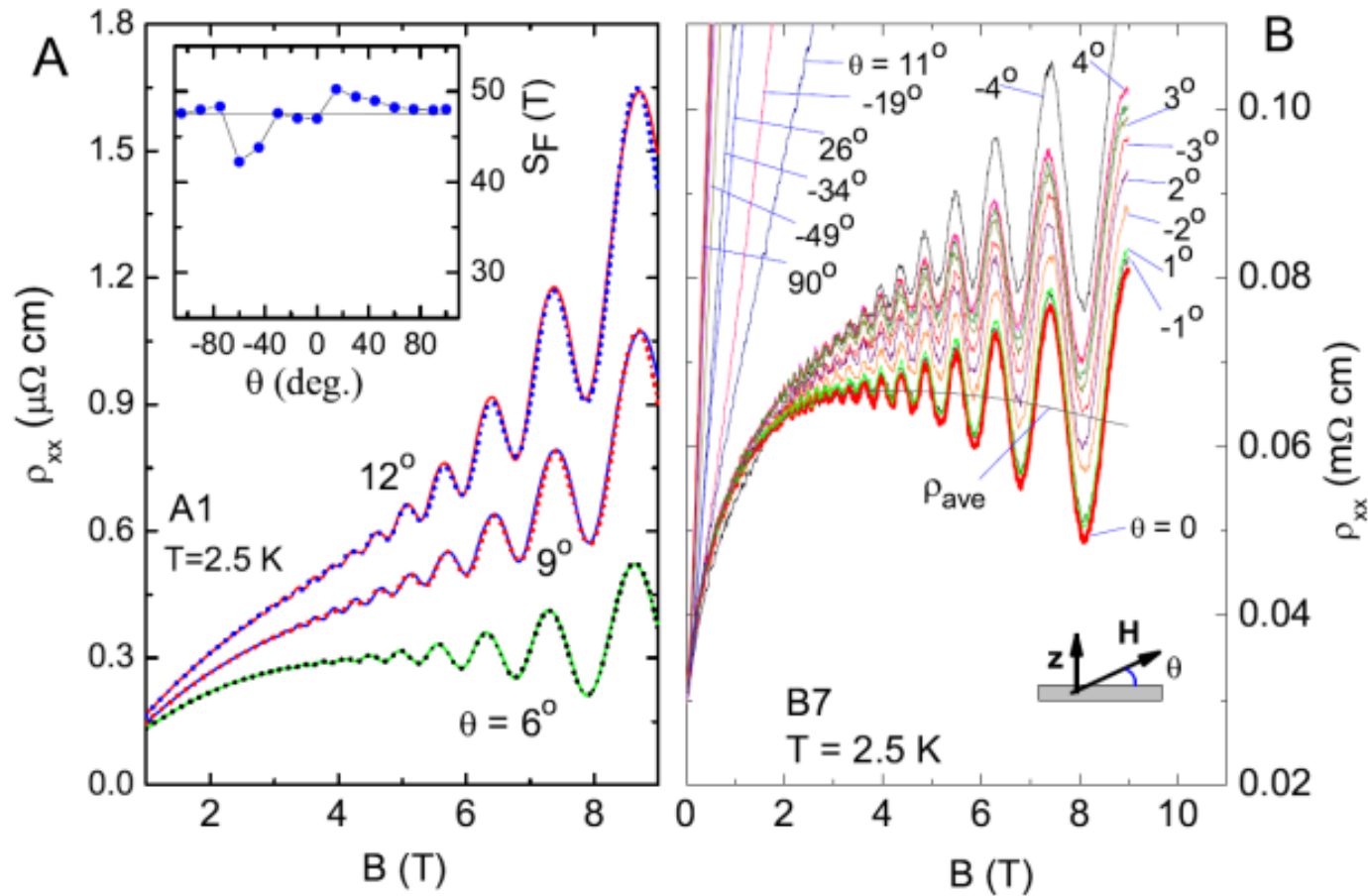
Some crystals show very low residual resistivity 20 n Ω cm

Correlated with large resistivity anisotropy (30)

Set B samples have higher resistivity

Shubnikov de Haas oscillations (Set B)

Tian Liang, Quinn Gibson *et al.* Nat. Mater. 2015



Chiral anomaly?

Negative, longitudinal MR detected at $\theta = 0^\circ$, but swamped by positive MR term (E_F is much too high).

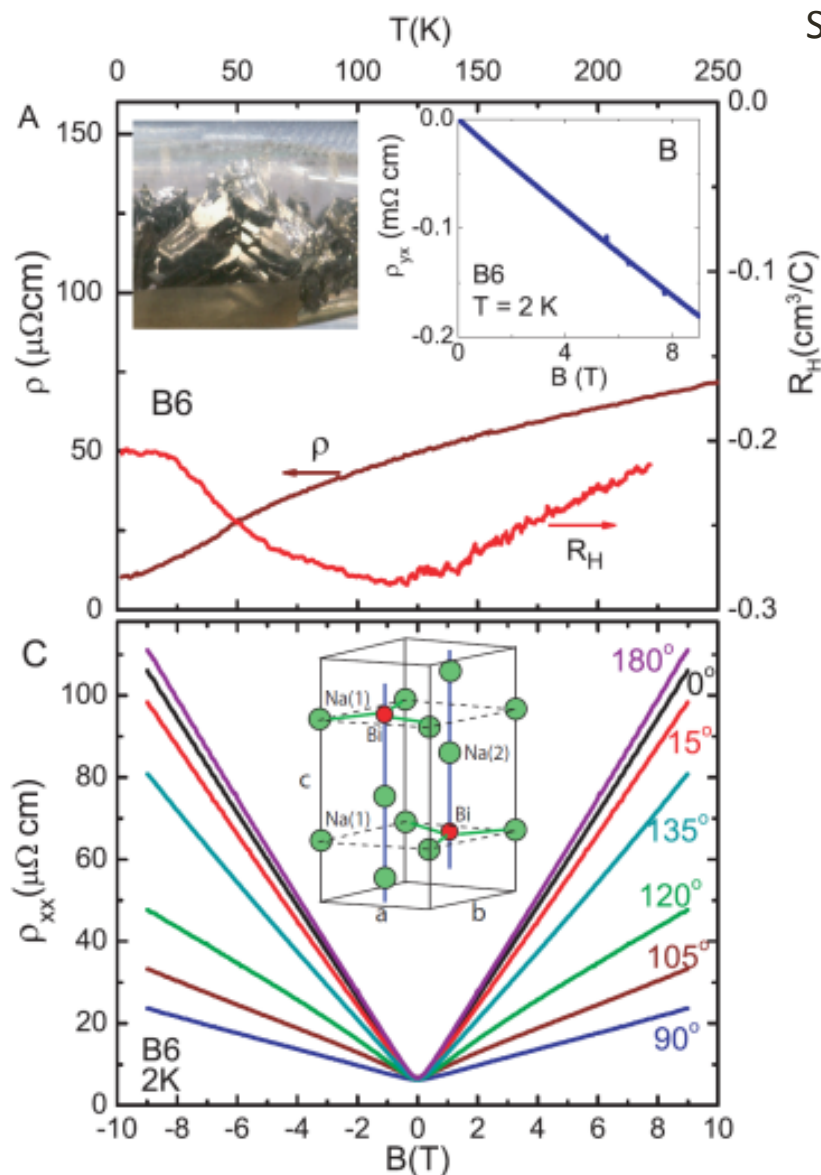
Transport results on Na₃Bi

S. Kushwaha *et al.*, *APL* 2015

Jun Xiong *et al.*, submitd



Jun Xiong Kushwaha Krizan



Deep purple crystals

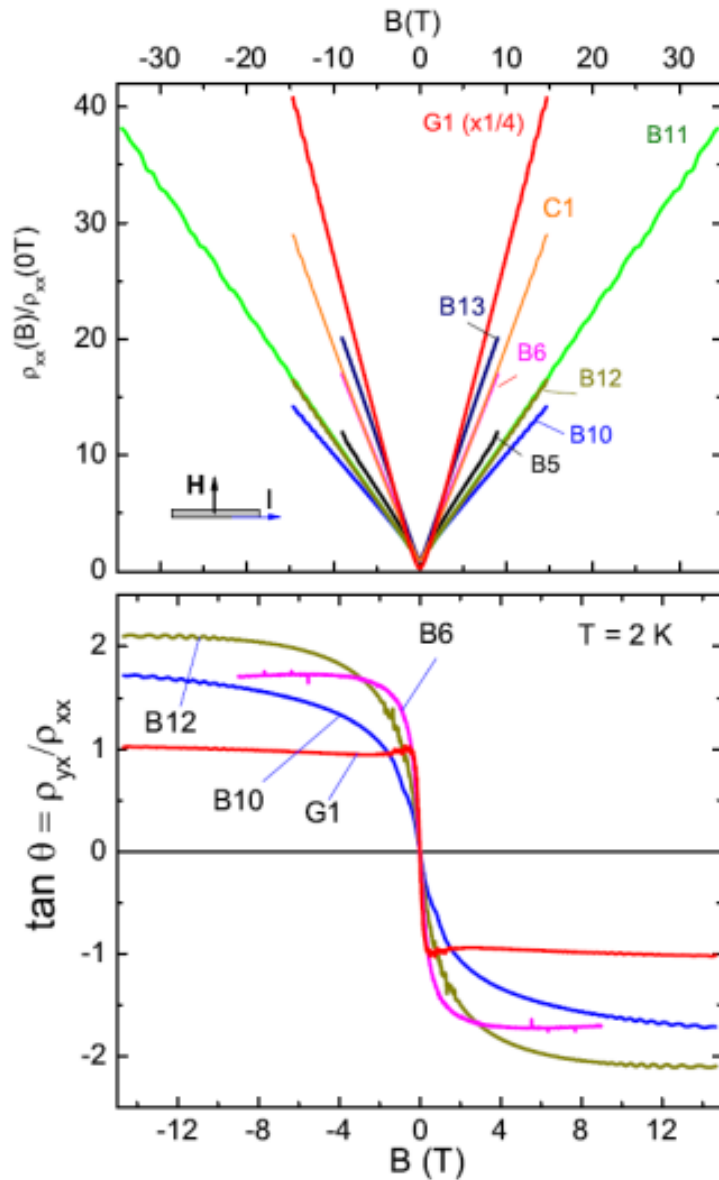
Rapidly oxidized in
ambient air (30 s)

Large linear MR similar to
Set B Cd₃As₂ samples

E_F 400 mV above node

H-linear MR and step-profile of Hall angle $\tan \theta_H$

Jun Xiong et al., PRL submitd



Conventional

$$\rho_{xx}(B) \sim [1 + (\mu B)^2]$$

$$\tan \theta_H \sim \mu B$$

In Na_3Bi

$$\rho_{xx}(B) \sim B, \quad \text{linear MR}$$

$\tan \theta_H$ has a *step-function* profile

Breaking Time Reversal symmetry in magnetic field

Search for Weyl physics

Chiral anomaly and axial current

Recent reports of Negative Longit MR in semimetals

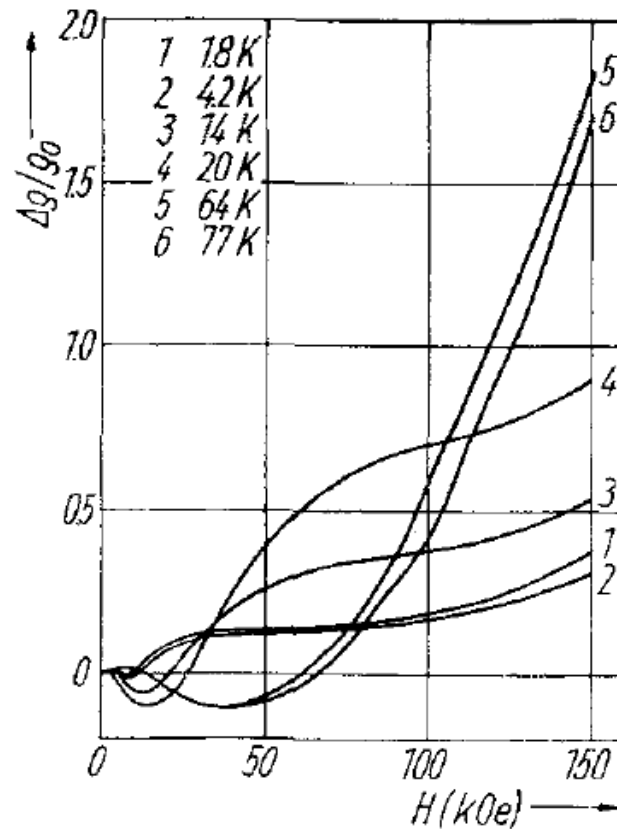
1. Neg. longit. MR in $\text{Bi}_{1-x}\text{Sb}_x$, Li et al., PRL 2013 --- **Accidental band crossing**
2. Neg. longit MR in Cd_3As_2 , Tian Liang et al., Nat Mat. 2015 --- **Dirac SM**
3. Neg. longit. MR in ZrTe_5 , Brookhaven Nat. Lab. arXiv --- **QSHE?**
4. Neg. longit. MR in PdCoO_2 , Balicas, Maeno, Hussey et al. arXiv --- **??**
5. Neg. longit. MR in TaAs, NbAs, S. Jia and Hasan, arXiv -- **Weyl metals**

A very old system

6. $\text{Cd}_x\text{Hg}_{1-x}\text{Te}$, Tsidilkovskii et al., Phys. Stat. Sol. (1974)

A very early report of negative longitud. MR in CdHgTe

Tsdidilkovskii, Girit, Kharus, Neifeld, Phys. Stat. Sol. 64, 717 (1974)



A negative longitud. MR is not restricted to chiral anomaly

Chiral anomaly in Quantum Field Theory

1969 Charged pions decay leptonically, $\pi^+ \rightarrow \mu^+ + \nu$

But, neutral pion decays electromagnetically

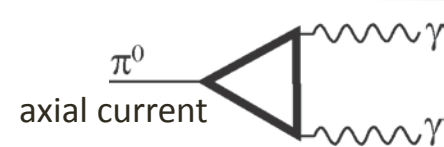
$$\pi^0 \rightarrow \gamma + \gamma$$

Chiral anomaly speeds up decay by a factor 3×10^8 !

Earliest (?) hint of 3 colors



Intro to QFT, Peskin Schroeder
QFT in a Nutshell, Tony Zee



Adler-Bell-Jackiw anomaly

Spoiler role in chiral gauge theories

Ruins chiral charge conservation and renormal.

Topological in origin

$$\nabla \cdot \mathbf{J}_5 + \frac{\partial \rho_5}{\partial t} = W$$

Anomaly-free condition

In a chiral theory, *all* the anomalies must cancel for it to be renormalizable.

This imposes stringent constraints, e.g. number of generations.

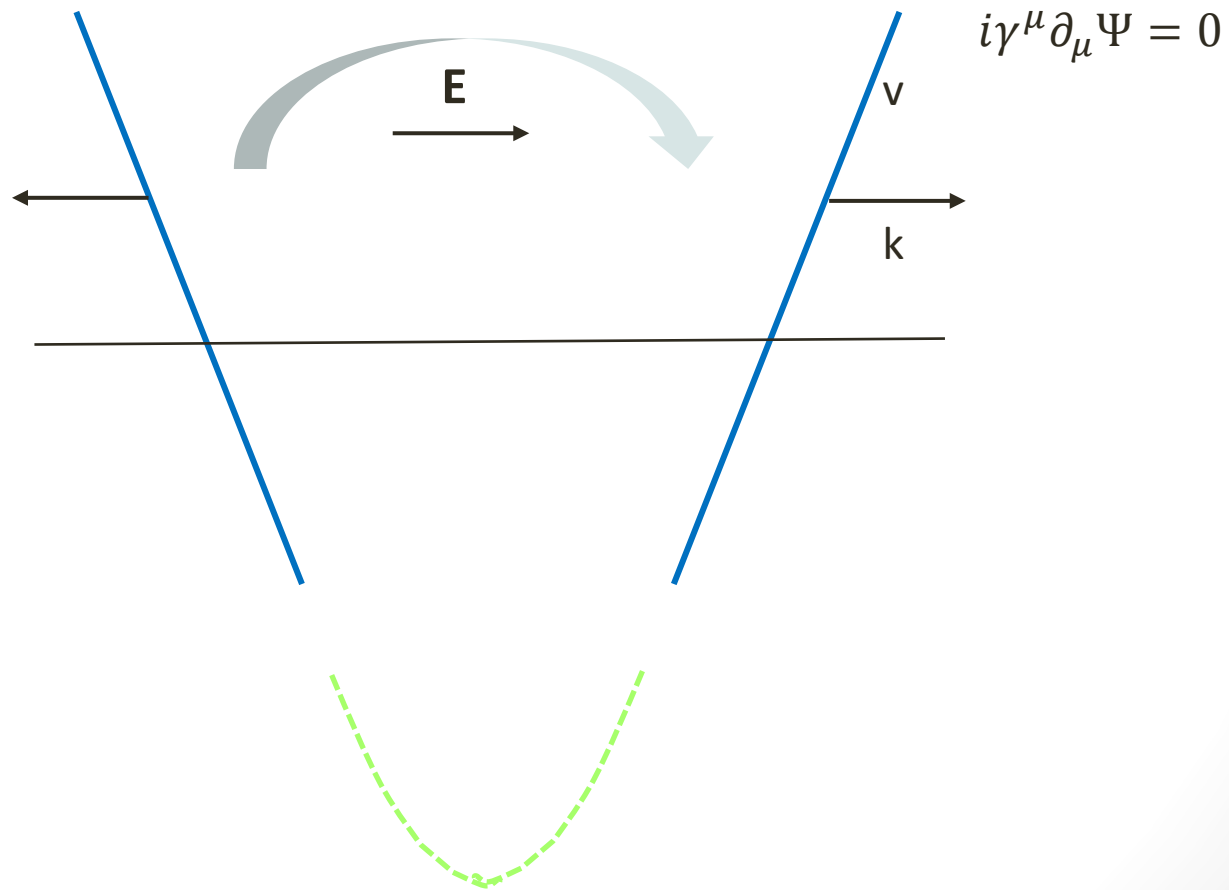
In Glashow-Weinberg-Salaam theory, exact cancellation involving 10+ diagrams has been called “magical”.

Nielsen and Ninomiya (*Phys. Lett.* 1983) proposed chiral anomaly may be observable in $(1+1)d$ or $(3+1)d$ crystals ---- even spacetime dim.

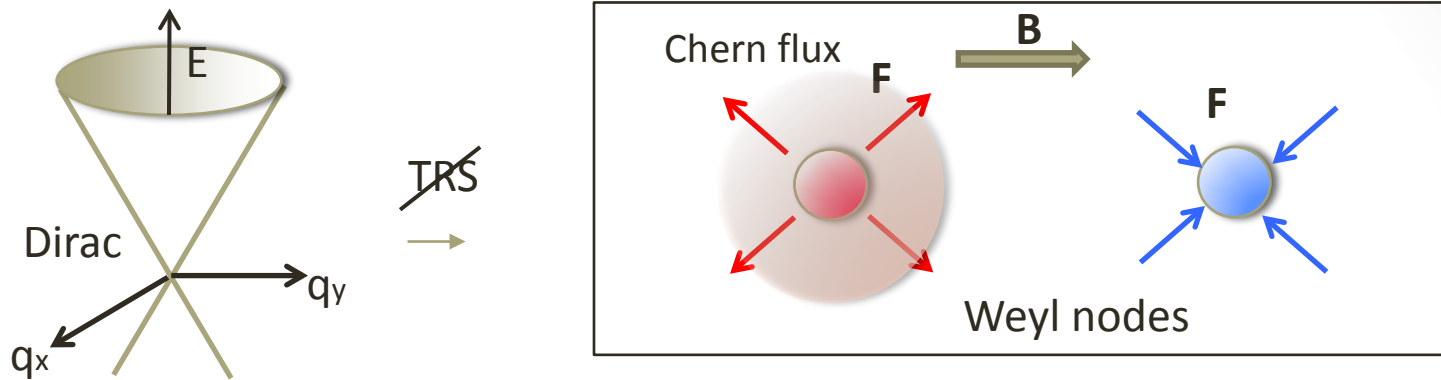
Detecting the chiral anomaly in a crystal?

Nielsen and Ninomiya (*Phys. Lett.* 1983)

For massless Dirac fields, the natural repres. is chiral (Weyl) representation. However, coupling to EM fields breaks chiral symmetry.



TRS breaking splits Dirac node into 2 Weyl nodes



Berry curvature $\mathbf{F}(\mathbf{k}) = \nabla_{\mathbf{k}} \times \mathbf{A}(\mathbf{k})$, $\mathbf{A}(\mathbf{k}) = i(u_{\mathbf{k}} | \nabla_{\mathbf{k}} u_{\mathbf{k}})$

- Weyl nodes come in pairs with $\chi = \pm 1$
- Acts as monopole source or sink of Berry curvature \mathbf{F} (or Chern flux Φ)

$$\chi = \frac{1}{2\pi} \oint \mathbf{F}(\mathbf{k}) \cdot d\mathbf{S}(\mathbf{k})$$

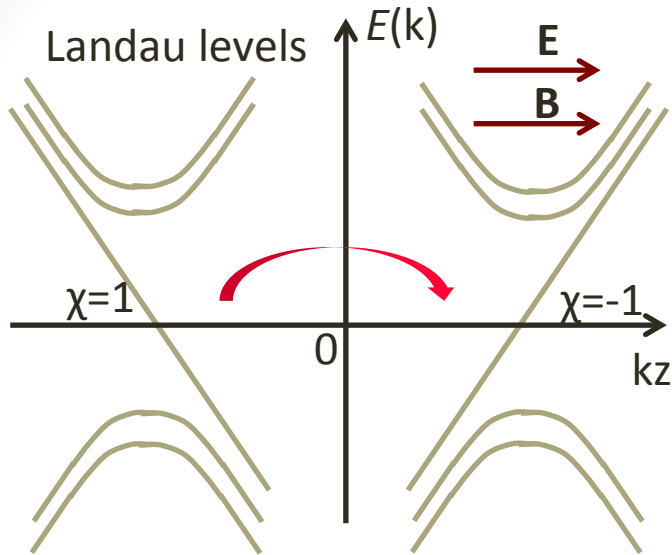
Non-conservation of chiral charge in E field \rightarrow the chiral anomaly

Wan, Turner, Vishwanath, *PRB* 2011

Burkov, Hook, Balents, *PRB* 2011 ... and 80+ theory uploads on arXiv

Son, Spivak, *PRB* 2013

Charge pumping and the chiral anomaly



Nielsen, Ninomiya, *Phys. Lett.* 1983
 Wan, Turner, Vishwanath, *PRB* 2011
 Burkov, Hook Balents, *PRB* 2011
 Son, Spivak, *PRB* 2013
 Parameswaran et al. *PRX* 2014
 Hosur and Qi, *Comp. Rend. Phys.* 2013

**Chiral anomaly engenders
 large, negative longitudinal MR
 Locked to B field**

In large- B regime, with $\mathbf{E} \parallel \mathbf{B}$, charge is pumped between Weyl nodes at the rate

$$W = -\frac{L^2}{2\pi\ell_B^2} \frac{Le\hbar}{2\pi} = -V \frac{e^3}{4\pi^2\hbar^2} \mathbf{E} \cdot \mathbf{B}$$

In weak B , charge pumping gives (Son and Spivak, *PRB* 2013)

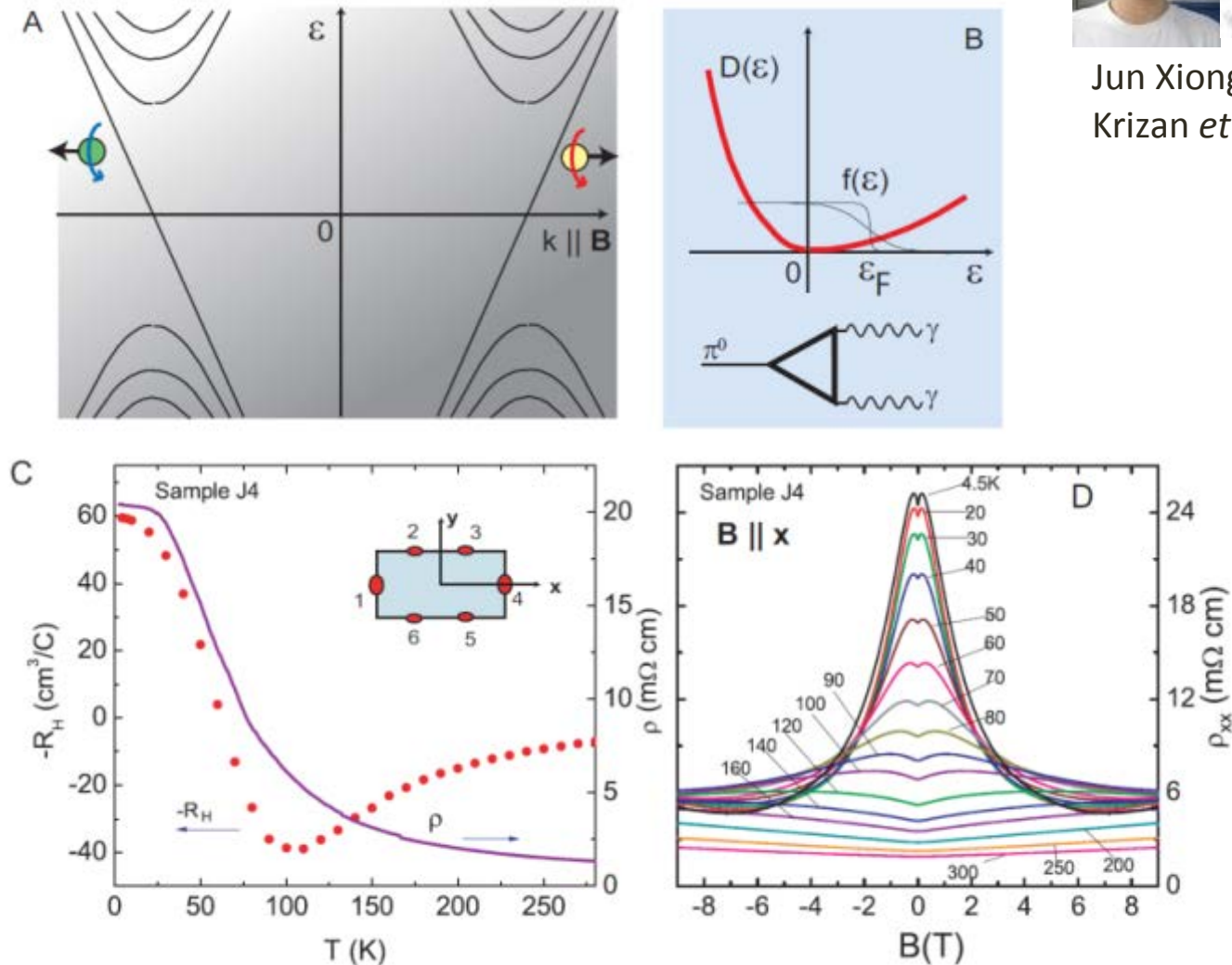
$$\sigma_\chi = \frac{e^2}{4\pi^2\hbar c} \frac{v}{c} \frac{(eBv)^2}{\epsilon_F^2} \tau_v, \quad \tau_v \text{ is intervalley scattg time}$$

Non-metallic resistivity and negative longitudinal MR in Na₃Bi

Long-term annealed crystals with $E_F \sim 30$ mV

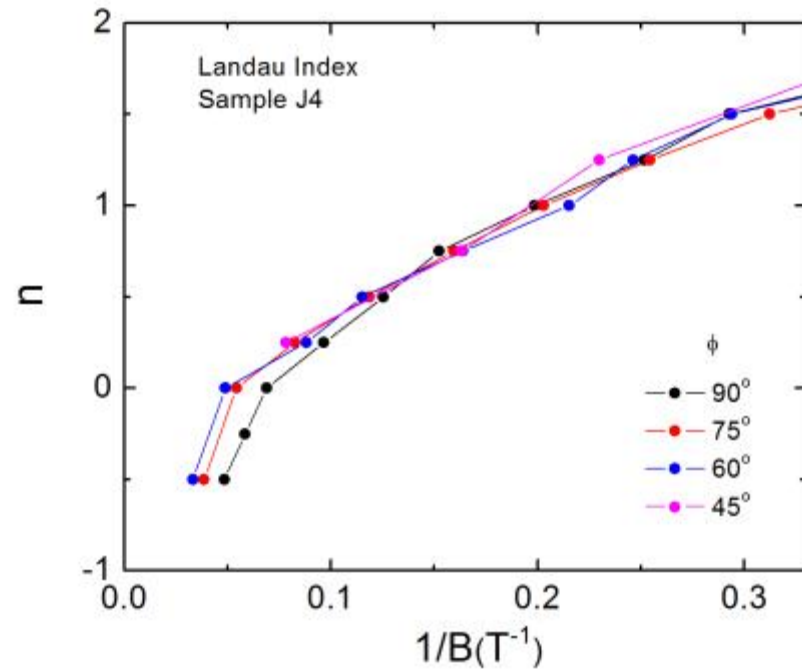
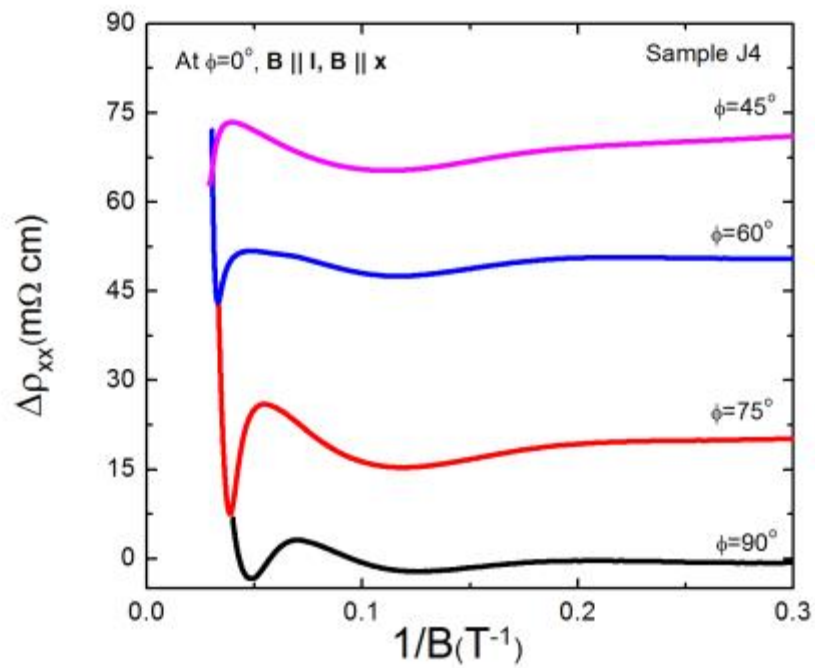


Jun Xiong, S. Kushwaha,
Krizan *et al.*, submitted

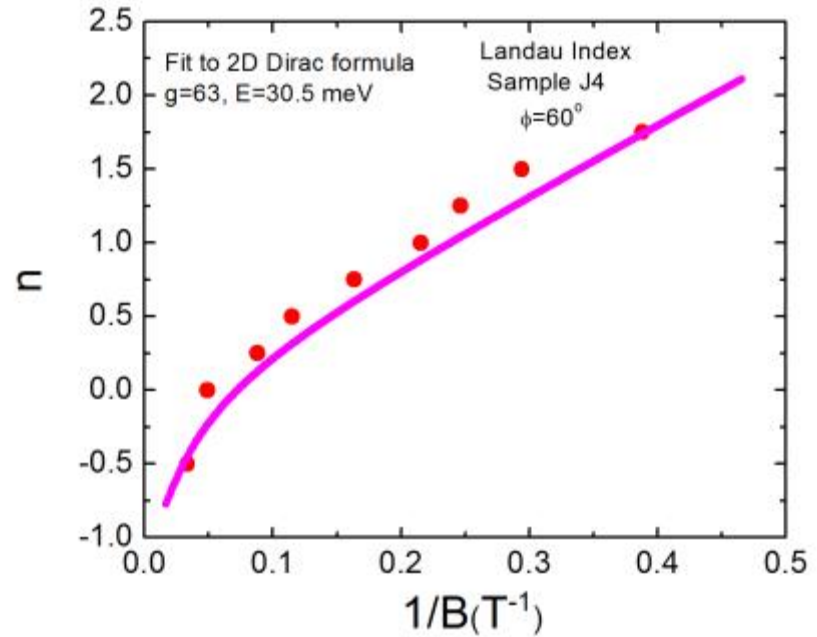
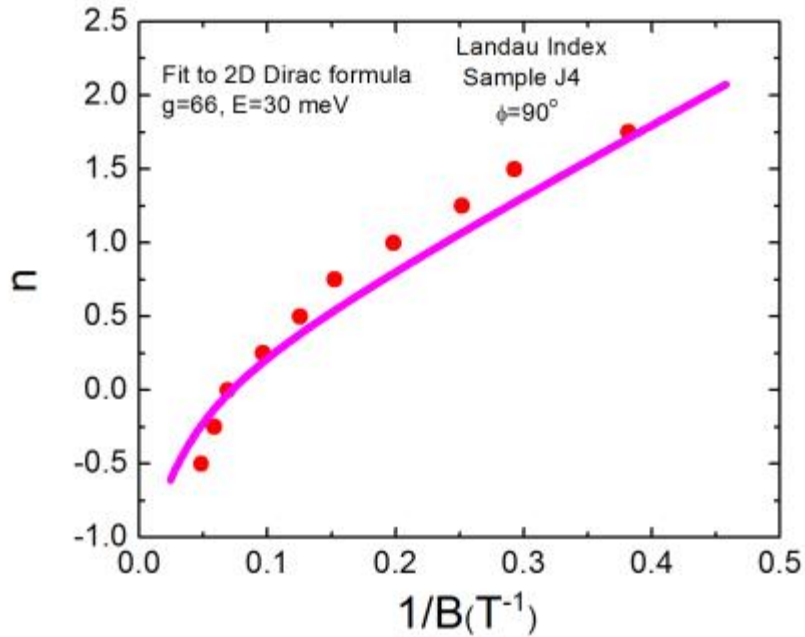


Extremum of RH implies a Fermi energy of 30 meV above node

Approaching $n=0$ Landau level



Evidence for a very large g factor (~ 63)

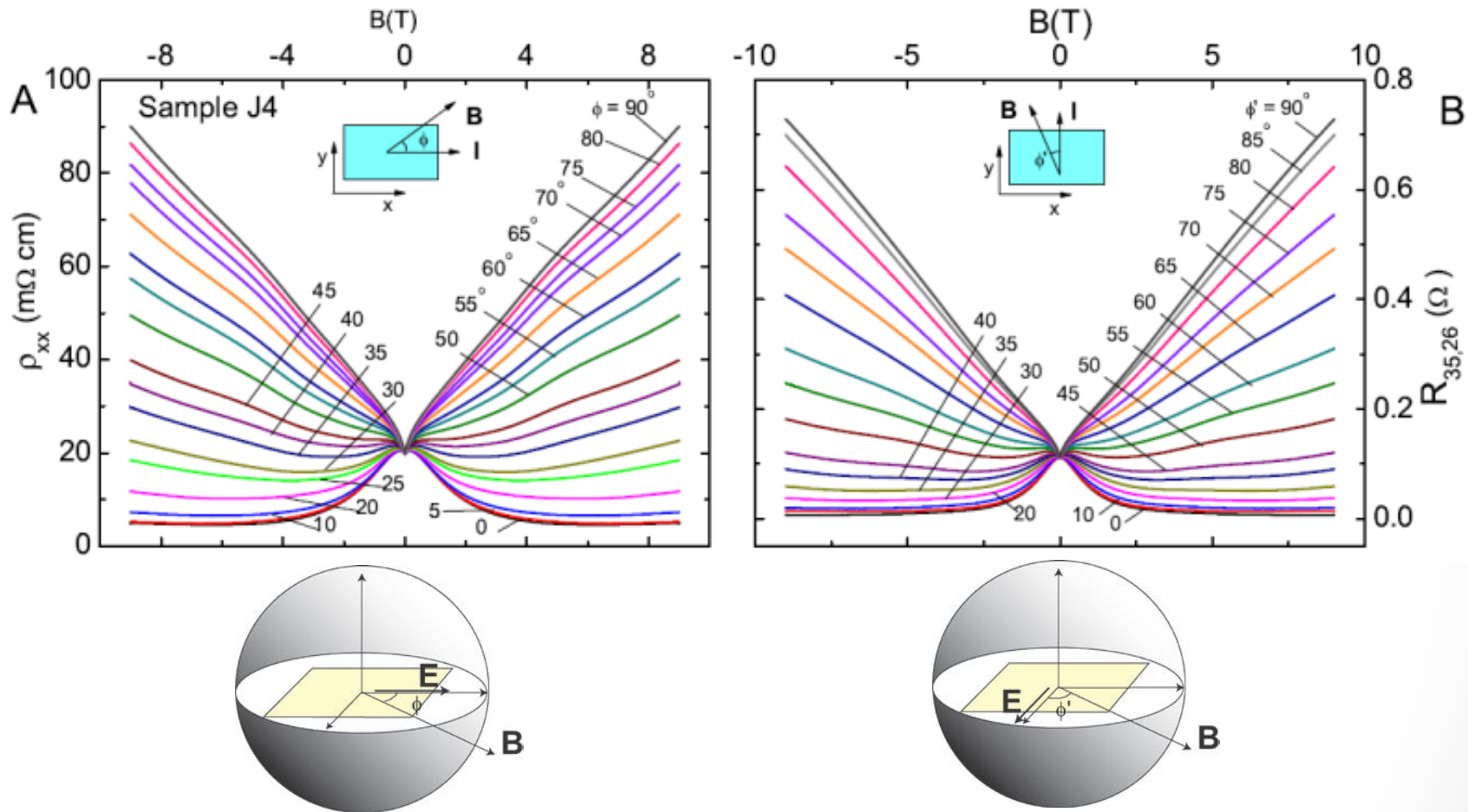


A very large deviation from straight line observed in index plot

Poor-man's fit yields $E_F = 30$ mV, $g \sim 63$

A crucial test for chiral anomaly -- **B** is locked to **E**

Jun Xiong, S. Kushwaha et al., submitted



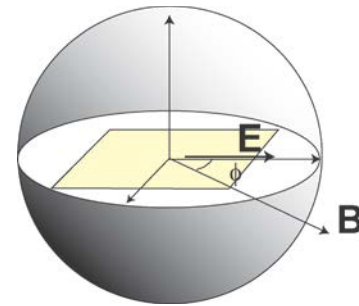
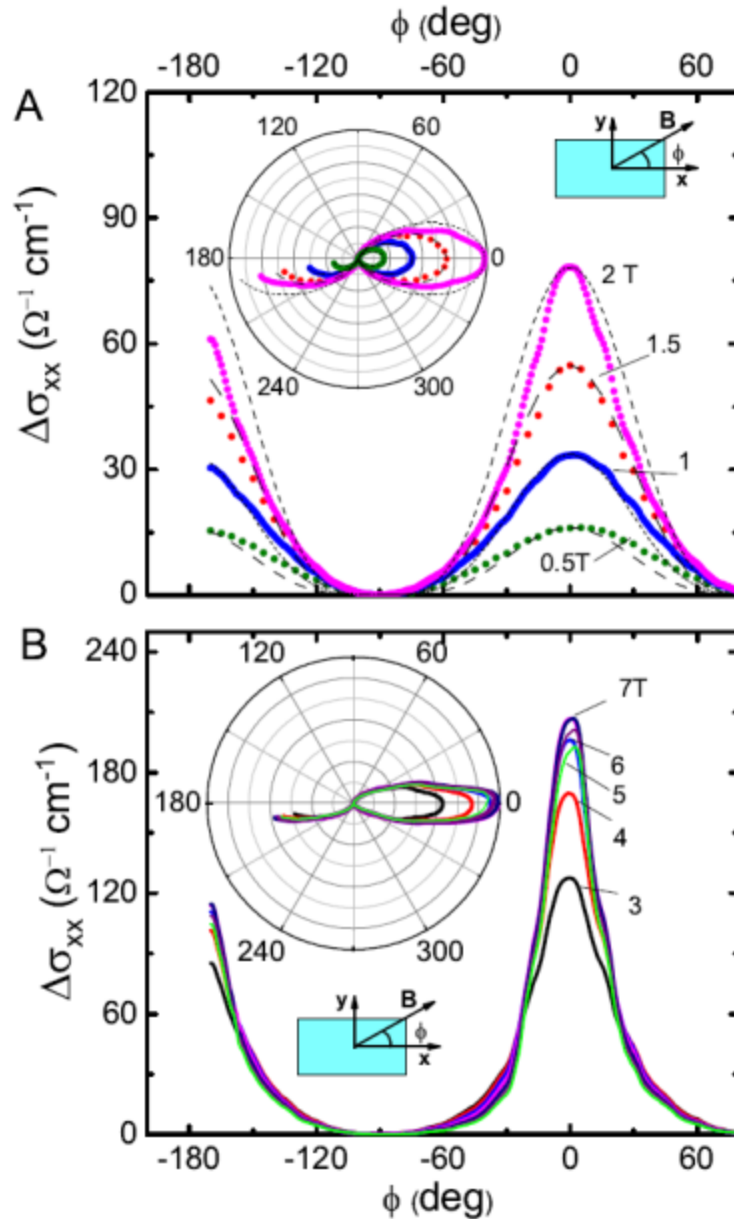
Negative MR appears only when **B** is locked to **E**.

Test: if **E** is rotated by 90° (right panel), neg. MR shifts to new direction of **E**.

For weak B, this locking is novel and unexpected in semiclassical transport

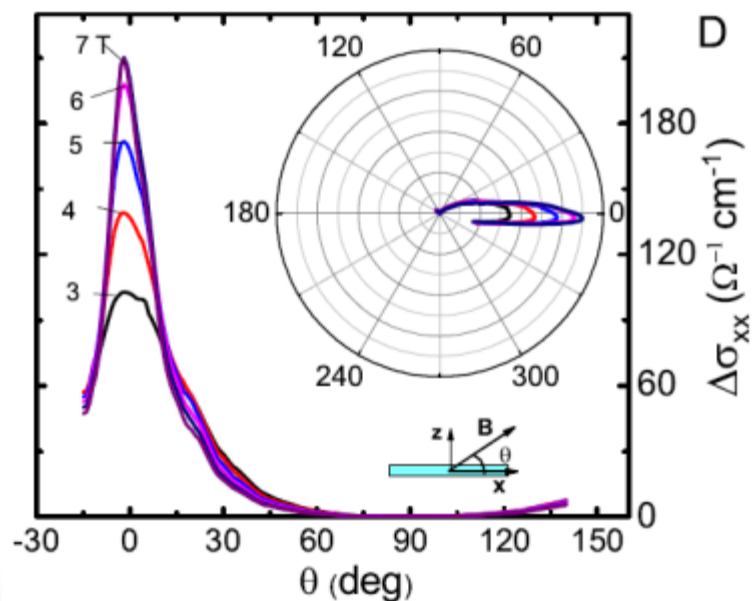
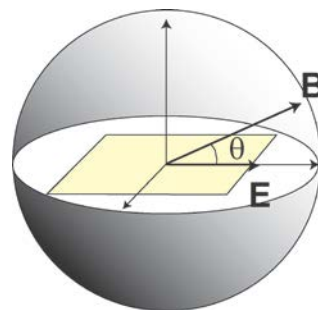
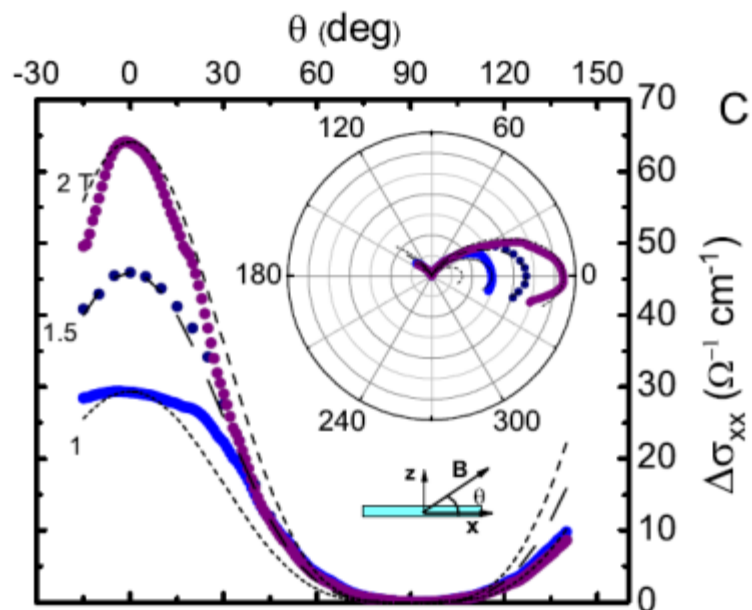
A narrow plume of chiral current, **B** in-plane

Jun Xiong, S. Kushwaha et al., submitted



Enhanced cond. in a narrowly collimated beam for **B** in the x-y (horizontal) plane

Width of chiral conductivity “plume”, \mathbf{B} normal to plane



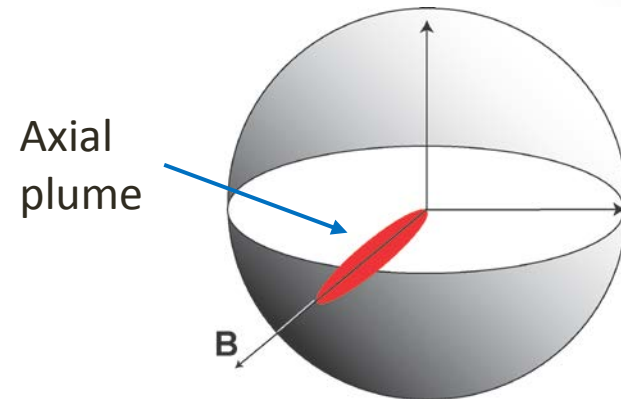
Enhanced cond. in a narrowly collimated beam for \mathbf{B} rotated in the x - z (vertical) plane

Signature of chiral anomaly: **B** locks direction of axial current

In conventional transport, we cannot “rotate” FS parameters, e.g. scattering rate anisotropy, by rotating direction of a weak **B**.

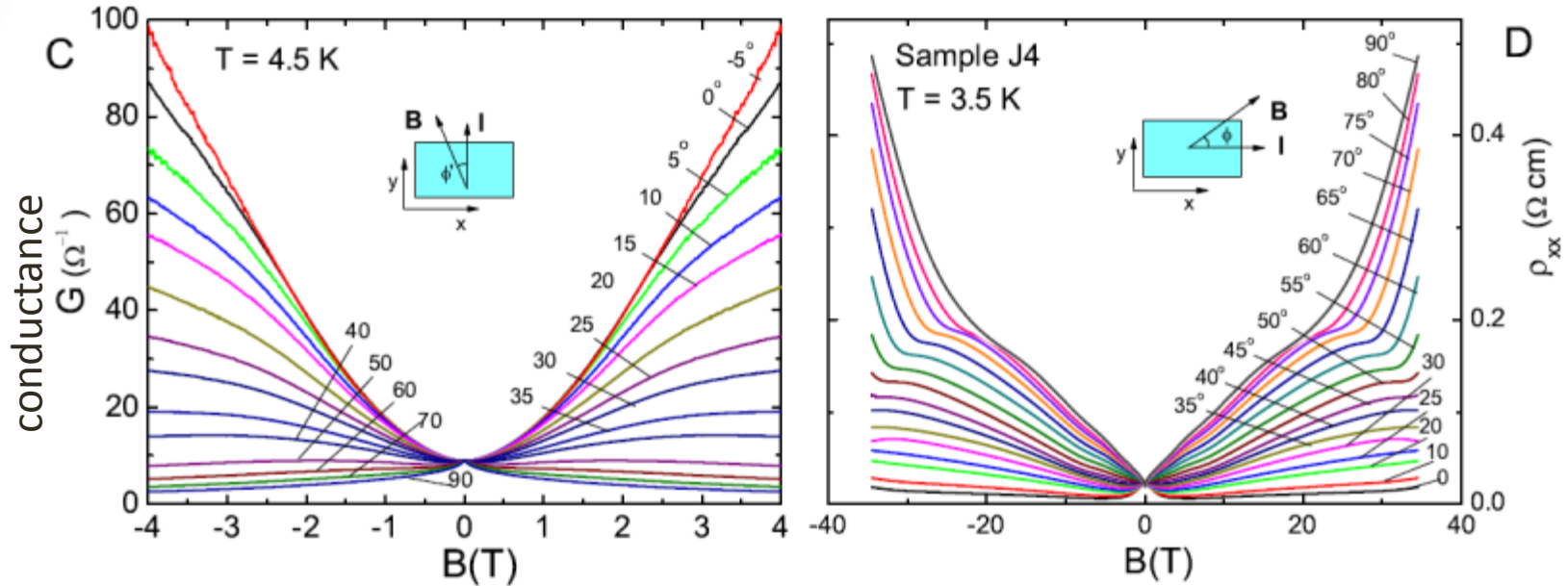
Locking of observed axial to **B** (even in *weak B*) seems to be a signature characteristic of the chiral anomaly.

Axial plume direction fixed by **B** (and **E**)



Observation of negative, longitudinal MR is necessary but insufficient.

Determine intervalley scat. lifetime using Son-Spivak expression



$$\sigma_x = \frac{e^2}{4\pi^2 \hbar c} \frac{v (eBv)^2}{\epsilon_F^2} \tau_v,$$

In semiclassical regime, Weyl node conductivity grows as B^2 (Son, Spivak, *PRB* '13)

From fit, we find $\tau_v = 40\text{-}80 \tau_0$

Summary

Transport experiments on Dirac semimetals

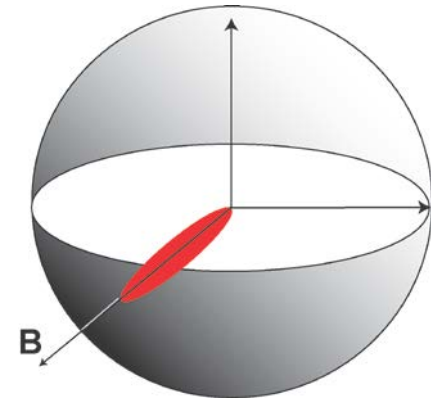
In Cd_3As_2 , we observe very high mobility (9 million cm^2/Vs) in zero B . Appears to be protected by a zero- B mechanism. Giant MR observed when protection is lifted.

In Na_3Bi , in samples with $E_F \sim 30$ mV, we see evidence for the chiral anomaly.

Signature: The enhanced-conductivity “plume” is locked to the direction of \mathbf{B} (and \mathbf{E}). **YES**

Estimated inter-valley lifetime is 40-80 x longer than Drude value.

A surprise: Width of plume is *much narrower* than anticipated by theory.



End



Jun Xiong



S. Kushwaha



Tian Liang



Jason Krizan



Quinn Gibson



Bob Cava



NPO

Thank you

References:

- 1) "Signature of the chiral anomaly in a Dirac semimetal: a current plume steered by a magnetic field," Jun Xiong, Satya K. Kushwaha, Tian Liang, Jason W. Krizan, Wudi Wang, R. J. Cava, N. P. Ong. [arXiv:1503.08179](https://arxiv.org/abs/1503.08179)
- 2) "Evidence for the chiral anomaly in a Dirac semimetal," Jun Xiong, Satya K. Kushwaha, Tian Liang, Jason W. Krizan, Max Hirschberger, Wudi Wang, R. J. Cava, N. P. Ong, submitted.